

Material descriptions

Feb. 10, 2014

- Absorption and dispersion in terms of susceptibility (Sect. 5.5):

Electric property measured by permittivity

$\epsilon = \epsilon_0(1 + \chi)$ and $\chi = \chi' + j\chi''$; real part relates to phase (dispersion) and imaginary part relates to amplitude (absorption) since propagation factor $\exp(-jkz)$ has $k = \omega\sqrt{\epsilon\mu_0}$

- Absorption (attenuation) coef (α) and refractive index n :

$$n = 1 + \chi'/2 \text{ and } \alpha = -k_0\chi''$$

Note: α and n are functions of freq.

- Measure of absorption and dispersion:

Attenuation coef in dB/km or m^{-1}

Describing freq dependent n -- group velocity

$$v_g = d\omega/d\beta = c_0/N, \text{ group index } N = n - \lambda_0 dn/d\lambda_0$$

where β is wave number.

$$\text{Material dispersion coef } D_\lambda = -\frac{\lambda_0}{c_0} \frac{d^2n}{d\lambda_0^2} \text{ (ps / km-nm).}$$

pulse widening or delay = $|D_\lambda|\Delta\lambda z$ (ps), where $\Delta\lambda$ is linewidth in nm, z is length of fiber in km.

- Kramers-Kronig relations:

Absorption and refractive index are connected by these relations; result of causality. (See Appendix B.1)

- Harmonic oscillator model for media:

Susceptibility is result of a sea of electric dipole

driven by an external electric field.

$d^2\vec{P}/dt^2 + \sigma d\vec{P}/dt + \omega_0^2\vec{P} = \omega_0^2\epsilon_0\chi_0\vec{E}$ where ω_0 is resonant freq, χ_0 is dc susceptibility, σ is damping coef

and $\vec{P} = \epsilon_0\chi\vec{E}$

$\chi = \chi_0 \frac{\nu_0^2}{\nu_0^2 - \nu^2 + j\nu\Delta\nu}$; $\Delta\nu = \sigma/2\pi$ is width of the resonance.

• Modeling dispersive system:

$A(z, t) = F^{-1}[H(f, z)F[A(0, t)]]$ where

$H(f) = \exp(-\alpha(f + \nu_0)z/2 - j[\beta(f + \nu_0) - \beta_0]z)$, ν_0 is the carrier (central) freq, f is the modulation freq, β_0 is

the wave number at carrier freq, $\beta = 2\pi\nu n/c_0$ and

$\nu = \nu_0 + f$

Note: Assuming $\nu_0 \gg f$

HW #2.5 due 2/21/14

1. Problem 5.4-1 (page 192 1st Ed, p. 195 2nd Ed)

2. exercise 3.1-4 (page 91 1st Ed, p. 84 2nd Ed)

3. Find the peak intensity and beam radius of a Gaussian beam by solving paraxial wave equation (Eq. (3.1-2) in the textbook) numerically. More precisely, you will solve the normalized paraxial

wave equation $-j \frac{\partial A}{\partial Z} + \nabla_{\perp}^2 A = 0$ where $Z = z/(4z_0)$, $(X, Y) = (x/w_0, y/w_0)$.

Recall the paraxial wave equation can be expressed in Fourier domain as:

$\frac{\partial A_F}{\partial Z} = j(K_x^2 + K_y^2)A_F$ where $A_F = \text{FFT}(A)$. Now the form of solution of A_F is just that of a 1st order ODE.

Consider a Gaussian beam with $\lambda_0 = 1\mu\text{m}$ propagating in air has $I_0 = 1\text{W}/\text{m}^2$ and $w_0 = 10\mu\text{m}$ at $z=0$. a)

Plot intensity distribution at $z=0$, b) plot intensity distribution at $z=2z_0$, c) find peak intensity and beam radius at $z=2z_0$. Your results should be obtained from Matlab and compared to analytical solutions.

Matlab function to be used: `fft2`, `ifft2`, `fftshift`, `image`, `colormap(gray(256))`. (Hint: to plot 8bit gray scale image, data must be normalized to 0 -255, i.e. input data array of image should be $255 * A / \max(A)$.)

4. Exercise 5.5-1 (page 176 in 1st Ed., page 172 in 2nd Ed.)

Extra-Credit

exercise 3.1-5 (page 91 1st Ed, p. 85 2nd Ed)